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# COMPOSITION OF SOME SINGULAR FOURIER INTEGRAL OPERATORS AND ESTIMATES FOR RESTRICTED $X$-RAY TRANSFORMS 

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## 0. Introduction.

Let $X$ and $Y$ be $C^{\infty}$ manifolds of dimension $n$ and $C \subset\left(T^{*} X \backslash 0\right) \times\left(T^{*} Y \backslash 0\right)$ a canonical relation; that is, $C$ is conic, smooth of dimension $2 n$ and the product symplectic form $\rho^{*} \omega_{X}-\pi^{*} \omega_{Y}$ vanishes identically on $T C$. (Here, $\omega_{X}, \omega_{Y}$ are the canonical symplectic forms on $T^{*} X, \quad T^{*} Y$, respectively, and $\rho: T^{*} X \times T^{*} Y \rightarrow T^{*} X$, $\pi: T^{*} X \times T^{*} Y \rightarrow T^{*} Y$ are the projections onto the first and second factors.) To $C$ is associated the class $I^{m}(C ; X, Y)$ of Fourier integral operators (FIOs) of order $m$ from $\mathscr{E}^{\prime}(Y)$ to $\mathscr{D}^{\prime}(X)$ ([18].) Composition calculi and sharp $L^{2}$ estimates for FIOs are only known under certain geometric conditions on the canonical relation(s). Most importantly, the transverse intersection calculus of Hörmander [18] implies that if $A_{1} \in I^{m_{1}}\left(C_{1} ; X, Y\right), A_{2} \in I^{m_{2}}\left(C_{2} ; Z, X\right)$ with $C_{1}$ and $C_{2}$ local canonical graphs, then $A_{2} A_{1} \in I^{m_{1}+m_{2}}\left(C_{2} \circ C_{1} ; Z, Y\right)$. In particular, if $C_{1}$ is a canonical graph, $A_{1}^{*} A_{1} \in I^{2 m_{1}}\left(\Delta_{T * Y} ; Y, Y\right)$ is a pseudodifferential operator and thus $A_{1}: H_{\text {comp }}^{s}(Y) \rightarrow H_{\text {loc }}^{s-m_{1}}(X)$ continuously, $\forall s \in \mathbf{R}$. Later, this composition calculus was extended by Duistermaat and Guillemin [9] and Weinstein [32] to the case of clean intersection.

[^0]For $L^{2}$ estimates, the following more general result holds ([18]). If the differentials of the mappings $\rho$ and $\pi$ drop rank by at most $k$, for some $k<n$, there is an estimate with a loss of $k / 2$ derivatives: $A: H_{\text {comp }}^{s}(Y) \rightarrow H_{\text {loc }}^{s-m-\frac{k}{2}}(X)$. This can be refined in the following way ([19], p. 30). Since $C$ is a canonical relation, on $C$ we have a closed 2-form $\omega_{C}=\rho^{*} \omega_{X}=\pi^{*} \omega_{Y}$, which is nondegenerate (i.e ; symplectic) iff $C$ is a local canonical graph. If $r$ is the co-rank of $C\left(=2 n-\operatorname{rank} \omega_{C} \leqslant 2 k\right)$, then $A: H_{\text {comp }}^{s}(Y) \rightarrow H_{\text {ioc }}^{s-m-\frac{r}{4}}(X)$. These results are sharp in that there are examples, such as the case when $C \subset\left(T^{*} X \backslash 0\right) \times\left(T^{*} Y \backslash 0\right)$ is the flowout of a codimension $k$ involutive submanifold of $T^{*} Y \backslash 0$, where one cannot do better. For canonical relations $C$ for which $\pi$ and $\rho$ become singular in specific ways, however, one expects there to be a sharp value $0<s_{0}=s_{0}(C) \leqslant \frac{r}{4} \leqslant \frac{k}{2} \quad$ such that $A: H_{\text {comp }}^{s}(Y) \rightarrow$ $H_{\text {loc }}^{s-m-s_{0}}(X), \forall s \in \mathbf{R}$. A result of this nature is contained in the work of Melrose and Taylor [25] on folding canonical relations, for which $\pi$ and $\rho$ have at most Whitney folds, so that $k=1, r=2$ and $\omega_{C}$ is a folded symplectic form. Via canonical transformations of $T^{*} X \backslash 0$ and $T^{*} Y \backslash 0, C$ can be conjugated (microlocally) to a single normal form ; on the operator level, $A$ can be conjugated by elliptic FIOs to an Airy operator on $\mathbf{R}^{n}$, from which the sharp boundedness $A: H_{\text {comp }}^{s}(Y) \rightarrow$ $H_{\text {loc }}^{s-m-1 / 6}(X)$ can be read off.

The purpose of the present work is to establish a composition calculus and obtain sharp $L^{2}$ estimates, with a loss of $\frac{1}{4}$ derivative, for a somewhat more singular class of canonical relations, the fibered folding canonical relations (FFCRs), for which again $\pi$ is a Whitney fold and $\omega_{C}$ is a folded symplectic form but for which $\rho$ is a «blowdown» ( $\simeq$ polar coordinates in two variables). These canonical relations arise naturally in integral geometry and were described independently in Greenleaf and Uhlmann [12] and Guillemin [15]. A specific canonical relation of this type had already been analyzed in considerable detail by Melrose [23]. Related operators are in Boutet de Monvel [3]. An unfortunate feature of FFCRs is that they cannot be conjugated to a single normal form. There are already obstructions to a formal power series attempt to derive a normal form (cf. [12]). Alternatively, as shown in [15], the canonical involution of $T^{*} X \backslash \rho(L)$, where $L \subset C$ is the fold hypersurface for $\pi$, induced by the $2-1$ nature of $\pi$ near $L$, may or may not extend smoothly past $\rho(L)$. In any event, it is not
possible to give exactly a phase function $\phi$ that parametrizes a general FFCR. A somewhat remarkable fact is that this difficulty disappears when one composes an $A \in I^{m}(C ; X, Y)$ with its adjoint. Our main result is

Theorem 0.1. - Let $C \subset\left(T^{*} X \backslash 0\right) \times\left(T^{*} Y \backslash 0\right)$ be a nonradial fibered folding canonical relation and $A \in I^{m}(C ; X, Y), B \in I^{m^{\prime}}\left(C^{t} ; Y, X\right)$ properly supported Fourier integral operators. Then $B A \in I^{m+m^{\prime}, 0}\left(\Delta_{T^{*} Y}, \Lambda_{\pi(L)}\right)$.

Here, $\Delta_{T^{*} Y}$ is the diagonal of $\left(T^{*} Y \backslash 0\right) \times\left(T^{*} Y \backslash 0\right), \pi(L) \subset T^{*} Y \backslash 0$ is the image of the fold hypersurface and $\Lambda_{\pi(L)}$ its flowout, and $I^{p, \ell}(\Delta, \Lambda)$ is the space of oscillatory integrals ("pseudodifferential operators with singular symbols») associated to the intersecting Lagrangians $\Delta$ and $\Lambda$ by Melrose and Uhlmann [26] and Guillemin and Uhlmann [16]. Using the estimates for elements of $I^{p, \ell}(\Delta, \Lambda)$ given in Greenleaf and Uhlmann [13], we obtain

Corollary 0.2. - For $A$ as above, $A: H^{s}(Y) \rightarrow H_{\text {loc }}^{s-m-\frac{1}{4}}(X)$ continuously, $\forall s \in \mathbf{R}$.

It should be remarked that the composition $A B$ is of a completely different nature, with the absence of a normal form for $C$ introducing serious analytical difficulties; this is discussed in Guillemin [15].

A special case of the theorem and corollary was proved in [13] for the restricted $X$-ray transform. If $(M, g)$ is an $n$-dimensional riemannian manifold for which the space $\mathscr{M}$ of (oriented) geodesics is a smooth ( $2 n-2$ )-dimensional manifold (e.g., $\mathbf{R}^{n}$ with the standard metric or a sufficiently small ball in any riemannian manifold), then the $X$-ray transform $\mathscr{R}: \mathscr{E}^{\prime}(M) \rightarrow \mathscr{D}^{\prime}(\mathscr{M})$ is given by

$$
\begin{equation*}
\mathscr{R} f(\gamma)=\int_{\mathbf{R}} f(\gamma(s)) d s, \quad \gamma \in \mathscr{M} \tag{0.3}
\end{equation*}
$$

$\gamma(s)$ any unit-velocity parametrization of $\gamma$. In the absence of conjugate points, $\mathscr{R}$ is an FIO of order $-\frac{n}{4}$ associated with a canonical relation satisfying the Bolker condition [14] and so $\mathscr{R}: H_{\text {comp }}^{s}(M) \rightarrow H_{\text {loc }}^{s-\frac{1}{2}}(\mathscr{M})$, generalizing (locally) the result of Smith and Solmon [28] on $\mathbf{R}^{n}$. (See also Strichartz [30] for the case of hyperbolic space.) Following Gelfand, one is also interested in the restriction of $\mathscr{R} f$ to $n$-dimensional
submanifolds $\mathscr{C} \subset \mathscr{M}$ (geodesic complexes); denote $\left.\mathscr{R} f\right|_{\mathscr{C}}$ by $\mathscr{R}_{\mathscr{C}} f$. Of particular interest are those $\mathscr{C}$ 's which are admissible for reconstruction of $f$ from $\mathscr{R}_{\mathscr{G}} f$ in that they satisfy a generalization of Gelfand's criterion [11]; in [12] it was shown that, with appropriate curvature assumptions, for such a $\mathscr{C}, \mathscr{R}_{\mathscr{C}}$ is an FIO of order $-\frac{1}{2}$ associated with a FFCR. In this case the Schwartz kernel of $\mathscr{R}_{\mathscr{8}}^{*} \mathscr{R}_{\mathscr{C}}$ is quite explicit and was shown in [13] to belong to $I^{-1,0}\left(\Delta_{T * M}, \Lambda_{\pi(L)}\right)$, yielding the boundedness of $\mathscr{R}_{\mathscr{C}}: H_{\text {comp }}^{s}(M) \rightarrow H_{\text {loc }}^{s+\frac{1}{4}}(\mathscr{C}), s \geqslant-\frac{1}{4}$.

To prove local $L^{p}$ estimates for admissible geodesic complexes, we extend $\mathscr{R}_{8}$ to an analytic family $R^{\alpha} \in I^{-\operatorname{Re}(\alpha)-\frac{1}{2}}(C ; \mathscr{C}, M)$; application of analytic interpolation then requires $L^{2}$ estimates for general elements of $I(C ; \mathscr{C}, M)$, for which the argument of [13] is insufficient. We prove

Theorem 0.4. - Let $\mathscr{C} \subset \mathscr{M}$ be an admissible geodesic complex and let $P(x, D)$ be a zeroth order pseudodifferential operator on $M$ such that $\mathscr{R}_{\mathscr{G}} P \in I(C ; \mathscr{C}, M)$ with $C$ a fibered folding canonical relation. Then $\mathscr{R}_{\mathscr{E}} P: L_{\text {comp }}^{p}(M) \rightarrow L_{\text {loc }}^{q}(\mathscr{M})$ for $p, q$ satisfying either of the following conditions:
(a) $1<p \leqslant \frac{4 n-3}{2 n-1}, \frac{1}{q} \geqslant \frac{2 n+1}{2 n p}-\frac{1}{2 n}$;
(b) $\frac{4 n-3}{2 n-1} \leqslant p<\infty, \frac{1}{q} \geqslant \frac{2 n-1}{2 n p}$.

For the full $X$-ray transform in $\mathbf{R}^{n}$, global $L^{p}$ estimates have been proven by Drury [6] [7] and refined by Christ [5] to mixed $L^{p}-L^{q}$ norms (see also [30], Oberlin and Stein [27]); however, even in $\mathbf{R}^{n}$ our estimates do not seem to be retrievable from theirs because of the high codimension of $\mathscr{C}$ in $\mathscr{M}$. Wang [31], using variations of the techniques of [5] [6] [7], has established global $L^{p}$ estimates for some special line complexes in $\mathbf{R}^{n}$.

There is a gap between the estimates in (0.4) and the expected optimal ones. Furthermore, one expects that, just as for the $L^{2}$ estimates [13], for general (nonadmissible) $\mathscr{C} \subset \mathscr{M}$, better estimates hold, reflecting the more singular way in which $C$ sits in $T^{*} \mathscr{C} \times T^{*} M$ when $\mathscr{C}$ is admissible. This is confirmed below for a particularly nice class of inadmissible $\mathscr{C}$ 's, for which $C$ is a folding canonical relation.

The paper is organized as follows. In §'1 we give a precise definition of FFCRs and recall the symplectic geometry needed to conjugate a FFCR into a position where it has a generating function $S\left(x, y_{n}, \eta^{\prime}\right)$. The geometry of $C$ then allows us to put a $S$ in a weak normal form. The relevant facts concerning $I^{p, \ell}(\Delta, \Lambda)$, including the iterated regularity characterization given in [13], are recalled in § 2 . In § 3 we prove (0.1) by computing $B A$, simplifying the phase, and then applying first order pseudodifferential operators to verify the iterated regularity condition. The applications to the restricted $X$-ray transform are given in $\S 4$.

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## 1. Weak normal form and phase functions.

Consider on $\mathbb{R}^{n} \times \mathbb{R}^{n} \times\left(\mathbb{R}^{n-1} \backslash 0\right)$ the phase function

$$
\begin{equation*}
\phi_{0}\left(x, y, \theta^{\prime}\right)=\left(x^{\prime}-y^{\prime}\right) \cdot \theta^{\prime}+\frac{x_{n}^{2} y_{n}}{2} \theta_{1}, \quad\left|\theta_{1}\right| \geqslant c|\theta|, y_{n} \neq 0 \tag{1.1}
\end{equation*}
$$

where we write $x=\left(x^{\prime}, x_{n}\right)=\left(x_{1}, x^{\prime \prime}, x_{n}\right) \in \mathbb{R}^{n}$. Calculating the critical set $\left\{\left(x, y, \theta^{\prime}\right): d_{\theta}{ }^{\prime} \phi_{0}=0\right\}$ and computing the map

$$
\left(x, y, \phi^{\prime}\right) \rightarrow\left(x, d_{x} \phi_{0} ; y,-d_{y} \phi_{0}\right),
$$

we find that $\phi_{0}$ parametrizes the canonical relation

$$
\begin{align*}
& C_{0}=\left\{\left(y_{1}-\frac{x_{n}^{2} y_{n}}{2}, y^{\prime \prime}, x_{n}, \eta^{\prime}, x_{n} y_{n} \eta_{1} ; y, \eta^{\prime},-\frac{x_{n}^{2} \eta_{1}}{2}\right):\right.  \tag{1.2}\\
&=\left\{\left(x, \xi^{\prime}, x_{n} y_{n} \xi_{1} ; x_{1}+\frac{x_{n}^{2} y_{n}}{2}, x^{\prime \prime}, y_{n}, \xi^{\prime},-\frac{x_{n}^{2}}{2} \xi_{1}\right):\right. \\
&\left.\left(x, \xi^{\prime}, y_{n}\right) \in \mathbb{R}^{2 n},\left|\xi_{1}\right| \geqslant c\left|\xi^{\prime}\right|, y_{n} \neq 0\right\}
\end{align*}
$$

Denoting, as before, the projections $C_{0} \rightarrow T^{*} \mathbb{R}^{n} \backslash 0$ onto the first and second factors by $\rho$ and $\pi$, respectively, one sees immediately that $C_{0}$ is a local canonical graph away from $L=\left\{x_{n}=0\right\}$, where $\pi$ has a Whitney fold (defined below) ; $\pi(L)=\left\{\eta_{n}=0\right\} \subset T^{*} \mathbb{R}^{n} \backslash 0$ is an embedded
hypersurface. At $L, \rho$ is more singular: $\rho(L)=\left\{x_{n}=\xi_{n}=0\right\} \subset T^{*} \mathbb{R}^{n} \backslash 0$ is embedded, codimension 2 , and symplectic (i.e. $\left.\sum_{1}^{n} d \xi_{j} \wedge d x_{j}\right|_{\rho(L)}$ is nondegenerate), and $\rho$ «blows up» $\rho(L)$, having 1-dimensional fibers with tangents $\frac{\partial}{\partial y_{n}} \cdot C_{0}$ is an example of a fibered folding canonical relation ; we recall from [12] and [15] the general definition of a FFCR and then show that any such can be conjugated sufficiently close to $C_{0}$ so that it has a phase similar to $\phi_{0}$

Definition 1.3. - Let $M$ and $N$ be $n$-dimensional manifolds; $f$ : $M \rightarrow N C^{\infty}$.
a) $f$ is a Whitney fold if near each $m_{0} \in M, f$ is either a local diffeomorphism or df drops rank simply by 1 at $m_{0}$, so that $L=\{m \in M$ : $\operatorname{rank}(d f(m))=n-1\} \quad$ is $\quad a \quad$ smooth hypersurface through $m_{0}$, and $\operatorname{ker}\left(d f\left(m_{0}\right)\right) \notin T_{m_{0}} L$.
b) $f$ is a blow-down along a smooth hypersurface $K \subset M$ if $f$ is a local diffeomorphism away from $K$, while df drops rank simply by 1 at $K$, where Hess $f \equiv 0$ and $\operatorname{ker}(d f) \subset T K$, so that $\left.f\right|_{K}$ has 1-dimensional fibers ; furthermore, letting, for $m_{0} \in K$,

$$
\overline{d f}: f^{-1}\left(f\left(m_{0}\right)\right) \rightarrow G_{n-1, n}\left(T_{f\left(m_{0}\right)} N\right)
$$

be the map sending $m$ to the hyperplane $d f(m)\left(T_{m} M\right) \subset T_{f\left(m_{0}\right)} N$, we demand that $d(\overline{d f})(v) \neq 0, v \in \operatorname{ker}\left(d f\left(m_{0}\right)\right) \backslash 0$.

Remark. - In [12], a blow-down was called a fibered fold. Since this terminology is apparently not standard, we have dropped it.

Definition 1.4. - Let $X$ and $Y$ be n-dimensional $C^{\infty}$ manifolds and $C \subset\left(T^{*} X \backslash 0\right) \times\left(T^{*} Y \backslash 0\right)$ a canonical relation. $C$ is a (nonradial) fibered folding canonical relation if
a) $\pi: C \rightarrow T^{*} Y \backslash 0$ is a Whitney fold, with fold hypersurface $L$, and $\pi(L)$ an embedded nonradial hypersurface;
b) $\rho: C \rightarrow T^{*} X \backslash 0$ is a blow-down (necessarily along $L$ ), with $\rho(L)$ embedded, nonradial and symplectic, and $\rho: C \backslash L \rightarrow T^{*} X \backslash 0$ is $1-1$.

In [12], an additional compatibility condition was imposed; namely, that the fibers $\left.\rho\right|_{L}$ be the lifts by $\pi$ of the bicharacteristic curves of $\pi(L)$. It was shown by Guillemin [15] that this is automatically satisfied.

By suitable choice of coordinate systems, the projections $\pi$ and $\rho$ may each be put into normal form ; the lack of a normal form for FFCRs stems from the inability to reconcile these coordinate systems in general. We recall

Proposition 1.5 (Melrose, [20]). - Let $M$ and $N$ be conic manifolds of dimension $2 n$, with $N$ symplectic. Suppose $f: M \rightarrow N$ has a Whitney fold along $L \ni m_{0}$ and $f(L)$ is non radial at $f\left(m_{0}\right)$.

Then there exist canonical coordinates on $N$ near $f\left(m_{0}\right)$ and coordinates $(s, \sigma)$ near $m_{0}$ on $M$, homogeneous of degrees 0 and 1 , respectively, with $s_{j}\left(m_{0}\right)=\delta_{n j}, \sigma_{j}\left(m_{0}\right)=\delta_{1 j}, \forall j$, such that $f(s, \sigma)=\left(s, \sigma^{\prime},-\frac{\sigma_{n}^{2}}{2 \sigma_{1}}\right)$.

Proposition 1.6. - Let $M$ and $N$ be as above. Suppose $g: M \rightarrow N$ is a blow-down along $L \ni m_{0}$ and $g(L)$ is nonradial and symplectic near $g\left(m_{0}\right)$. Then there exist canonical coordinates on $N$ near $g\left(m_{0}\right)$ and coordinates $(t, \tau)$ near $m_{0}$ on $M$, homogeneous of degrees 0 and 1 , respectively, with $t_{j}\left(m_{0}\right)=0, \quad \tau_{j}\left(m_{0}\right)=\delta_{1 j}+\delta_{n j}, \quad \forall j, \quad$ such that $g(t, \tau)=\left(t, \tau^{\prime}, t_{n} \tau_{n}\right)$.

Proof. - Without the homogeneity, this is Theorem 4.5 of [12]; the proof there is easily adapted to the conic setting using the version of Darboux' theorem in [21].

Now let $C$ be a FFCR and apply (1.5), (1.6) to $f=\pi, g=\rho$, respectively, to obtain canonical coordinates on $T^{*} Y \backslash 0, T^{*} X \backslash 0$ and homogeneous coordinates $(s, \sigma),(t, \tau)$ near $c_{0} \in L \subset C$. Let

$$
T_{1}=s_{1}-\frac{\sigma_{n}^{2} s_{n}^{2}}{2 \sigma_{1}^{2}}, \quad T_{n}=\frac{\sigma_{n}}{\sigma_{1}} \quad \text { and } \quad S_{n}=\frac{\tau_{n}}{\tau_{1}}
$$

so that with respect to the homogeneous coordinate systems ( $T_{1}, T_{n}, s^{\prime \prime}, s_{n}, \sigma^{\prime}$ ) and ( $t, \tau^{\prime} S_{n}$ ) near $c_{0}$,

$$
\begin{gather*}
\pi\left(T_{1}, T_{n}, s^{\prime \prime}, s_{n}, \sigma^{\prime}\right)=\left(T_{1}+\frac{T_{n}^{2}}{2} s_{n}, s^{\prime \prime}, s_{n}, \sigma^{\prime},-\frac{T_{n}^{2}}{2} \sigma_{1}\right)  \tag{1.7}\\
\rho\left(t, \tau^{\prime}, S_{n}\right)=\left(t, \tau^{\prime}, S_{n} t_{n} \tau_{1}\right)  \tag{1.8}\\
\omega_{C}=d \sigma_{1} \wedge d T_{1}+d \sigma^{\prime \prime} \wedge d s^{\prime \prime}+T_{n}\left(s_{n} d \sigma_{1}+\sigma_{1} d s_{n}\right) \wedge d T_{n} \\
=d \tau^{\prime} \wedge d t^{\prime}+t_{n}\left(S_{n} d \tau_{1}+\tau_{1} d S_{n}\right) \wedge d t_{n} ;
\end{gather*}
$$

and

$$
\begin{equation*}
L=\left\{T_{n}=0\right\}=\left\{t_{n}=0\right\} . \tag{1.10}
\end{equation*}
$$

A function $f \in C^{\infty}(C)$ has a (singular) Hamiltonian vector field with respect to the folded symplectic form $\omega_{c}$, which expressed in the ( $T_{1}, T_{n}, s^{\prime \prime}, s_{n}, \sigma^{\prime}$ ) coordinates is

$$
\begin{align*}
H_{f}^{c}= & \left(\frac{\partial f}{\partial \sigma_{1}}-\frac{s_{n}}{\sigma_{1}} \frac{\partial f}{\partial s_{n}}\right) \frac{\partial}{\partial T_{1}}+\frac{1}{T_{n} \sigma_{1}} \frac{\partial f}{\partial s_{n}} \frac{\partial}{\partial T_{n}}  \tag{1.11}\\
& +\sum_{j=2}^{n-1} \frac{\partial f}{\partial \sigma_{j}} \frac{\partial}{\partial s_{j}}-\frac{\partial f}{\partial s_{j}} \frac{\partial}{\partial \sigma_{j}} \\
& +\left(\frac{s_{n}}{\sigma_{1}} \frac{\partial f}{\partial T_{1}}-\frac{1}{T_{n} \sigma_{1}} \frac{\partial f}{\partial T_{n}}\right) \frac{\partial}{\partial s_{n}}-\frac{\partial f}{\partial T_{1}} \frac{\partial}{\partial \sigma_{1}}
\end{align*}
$$

On $L,\left\{S_{n}=1\right\}$ has the form $\left\{s_{n}=1+F\left(T_{1}, s^{\prime \prime}, \sigma^{\prime}\right)\right\}$, so we let

$$
f\left(T_{1}, T_{n}, s^{\prime \prime}, s_{n}, \sigma^{\prime}\right)=-\sigma_{1} F\left(T_{1}, s^{\prime \prime}, \sigma^{\prime}\right) \frac{T_{n}^{2}}{2}
$$

Then there is a smooth function on $T^{*} Y \backslash 0$, which we denote by $\pi_{*} f$, such that $\pi^{*}\left(\pi_{*} f\right)=f$; of course, $H_{\pi * f}$ is a $C^{\infty}$ vector field on $T^{*} Y \backslash 0$, with $\chi_{\pi * f}=\exp \left(H_{\pi * f}\right)$ a canonical transformation. On the other hand, $H_{f}^{c}=F \frac{\partial}{\partial s_{n}}+O\left(T_{n}^{2}\right)$ and is $C^{\infty}$ by (1.11), and the $\omega_{c}$-morphism $\chi_{f}^{C}=\exp \left(H_{f}^{c}\right)$ is of the form

$$
\chi_{f}^{C}\left(T_{1}, T_{n}, s^{\prime \prime}, s_{n}, \sigma^{\prime}\right)=\left(T_{1}, T_{n}, s^{\prime \prime}, s_{n}+F\left(T_{1}, s^{\prime \prime}, \sigma^{\prime}\right), \sigma^{\prime}\right)+O\left(T_{n}^{2}\right)
$$

Changing variables on $C$ and $T^{*} Y \backslash 0$ simultaneously, we retain (1.7) and (1.9), but now have $\left\{T_{n}=s_{n}-1=0\right\}=\left\{t_{n}=S_{n}-1=0\right\}$ near $c_{0}$; denote this smooth $(2 n-2)$ - dimensional manifold by $L_{0}$ and let $i$ : $L_{0} \hookrightarrow C$ be the inclusion map. From (1.9), we have

$$
i^{*} \omega_{C}=d \sigma_{1} \wedge d T_{1}+d \sigma^{\prime \prime} \wedge d s^{\prime \prime}=d \tau^{\prime} \wedge d t^{\prime}
$$

By Darboux we can find a canonical transformation $\chi_{0}$ of $\mathbb{R}^{2 n-2}$ such that $\chi_{0}^{*}\left(T_{1}, s^{\prime \prime}, \sigma^{\prime}\right)=\left(t^{\prime}, \tau^{\prime}\right)$. Extending $\chi_{0}$ to be independent of $T_{n}$ and $s_{n}$, we obtain an $\omega_{c}$-morphism $\chi$ such that

$$
\begin{aligned}
\chi^{*}\left(T_{1}, s^{\prime \prime}, \sigma^{\prime}\right)=\left(t^{\prime}, \tau^{\prime}\right)+O\left(t_{n}\right)+O\left(S_{n}-1\right) & , \chi^{*} s_{n}=1 \\
& +a S_{n}+O\left(\left(S_{n}-1\right)^{2}\right)+O\left(t_{n}\right)
\end{aligned}
$$

and $\chi^{*} T_{n}=b t_{n}$ with $a \neq 0, b \neq 0$ near $c_{0}$. On the other hand, by simultaneously applying $\chi_{0}$ in the ( $y^{\prime}, \eta^{\prime}$ ) variables, we preserve (1.7). Thus, we have $\rho^{*}(x)=t, \pi^{*}\left(y_{n}\right)=s_{n}$ and $\pi^{*}\left(\eta^{\prime}\right)=\sigma^{\prime}$ forming local coordinates on $C$ near $c_{0}$; furthermore, $L=\left\{x_{n}=0\right\}$ in these coordinates, $\pi(L)=\left\{(y, \eta): \eta_{n}=0\right\}$ and $\rho(L)=\left\{(x, \xi): x_{n}=\xi_{n}=0\right\}$, and $d \rho^{*}\left(d \xi_{n}\right) \neq 0$.

Since $\left(x, y_{n}, \eta^{\prime}\right)$ form coordinates on $C$, there exists a generating function $S\left(x, y_{n}, \eta^{\prime}\right)$ for $C([18]): S$ is $C^{\infty}$, homogeneous of degree 1 in $\eta^{\prime}$, and

$$
\begin{equation*}
C=\left\{\left(x, d_{x} S ; d_{\eta}^{\prime} S, y_{n}, \eta^{\prime}, d_{y_{n}} S\right):\left(x, y_{n}, \eta^{\prime}\right) \in U\right\} \tag{1.12}
\end{equation*}
$$

near $c_{0}$, where $U$ is a conic neighborhood of $x=0, y_{n}=1, \eta^{\prime}=d y_{1}$, and $\phi\left(x, y, \eta^{\prime}\right)=S\left(x, y_{n}, \eta^{\prime}\right)-y^{\prime} \cdot \eta^{\prime}$ parametrizes $C$ near $c_{0}$. The fact that $C$ is a FFCR imposes several conditions on $S$, which we next derive.

That $\pi(L)=\left\{\eta_{n}=0\right\}$ implies that $\frac{\partial S}{\partial y_{n}}\left(x^{\prime}, 0, y_{n}, \eta^{\prime}\right)=0$, whence $\left.S\right|_{\left\{x_{n}=0\right\}}$ is independent of $y_{n}: S\left(x^{\prime}, 0, y_{n}, \eta^{\prime}\right)=S_{0}\left(x^{\prime}, \eta^{\prime}\right)$ for some smooth, homogeneous $S_{0}$. Since $\rho(L)=\left\{x_{n}=\xi_{n}=0\right\}$, we have $\frac{\partial S}{\partial x_{n}}\left(x^{\prime}, 0, y_{n}, \eta^{\prime}\right)=0$, so that

$$
\begin{equation*}
S\left(x, y_{n}, \eta^{\prime}\right)=S_{0}\left(x^{\prime}, \eta^{\prime}\right)+\frac{x_{n}^{2}}{2} S_{2}\left(x, y_{n}, \eta^{\prime}\right) \tag{1.13}
\end{equation*}
$$

where $S_{2}$ is smooth and homogeneous of degree 1 in $\eta^{\prime}$. The matrix representing $d \pi$ is

$$
d \pi=\left[\begin{array}{ccc}
d_{\eta^{\prime} x}^{2} S & d_{\eta^{\prime} y_{n}}^{2} S & d_{\eta^{\prime} \eta^{\prime}}^{2} S  \tag{1.14}\\
0 & 1 & 0 \\
0 & 0 & I_{n-1} \\
d_{y_{n} x}^{2} S & d_{y_{n} y_{n}}^{2} S & d_{y_{n} n^{\prime}}^{2} S
\end{array}\right]
$$

By the above comments, at $x_{n}=0$ the $y_{n}$-row and the $x_{n}$-column vanish; but since $\pi$ is a fold, $\left.d \pi\right|_{d x_{n}=0}$ has rank $2 n-1$, and thus $\operatorname{det}\left(d_{\eta^{\prime} x^{\prime}}^{2} S\right) \neq 0$ at $x_{n}=0$, i.e.,
(1.15) $S_{0}\left(x^{\prime}, \eta^{\prime}\right)$ is a nondegenerate generating function,
in $(n-1)$ variables. Also, $\operatorname{ker}(d \pi)=\mathbf{R} \frac{\partial}{\partial y_{n}}$ at $x_{n}=0$. Additionally,

$$
d \rho=\left[\begin{array}{cccc} 
& \circ & \circ &  \tag{1.16}\\
I_{n-1} & \vdots & \vdots & O \\
& \circ & \circ & \\
\circ \cdots \circ & 1 & 0 & \\
& \circ & & \\
d_{x^{\prime} x^{\prime}}^{2} S & \vdots & d_{x^{\prime} y_{n}}^{2} S & d_{x^{\prime} n^{\prime}}^{2} S \\
0 \cdots \circ & \circ & \\
\circ \cdots \circ & d_{x_{n} x_{n}}^{2} S & \circ & \circ
\end{array}\right]
$$

The nondegeneracy of $d_{x^{\prime} \eta^{\prime}}^{2} S$ yields (at $x_{n}=0$ )

$$
\begin{equation*}
\operatorname{Im}(d \rho)=\operatorname{span}\left\{\left\{\frac{\partial}{\partial x_{j}}\right\}_{j=1}^{n-1}, \frac{\partial}{\partial x_{n}}+\frac{\partial^{2} S}{\partial x_{n}^{2}} \frac{\partial}{\partial \xi_{n}},\left\{\frac{\partial}{\partial \xi_{j}}\right\}_{j=1}^{n-1}\right\} \tag{1.17}
\end{equation*}
$$

From $d \rho^{*}\left(d \xi_{n}\right) \neq 0$ it follows that

$$
\begin{equation*}
\frac{\partial^{2} S}{\partial x_{n}^{2}}\left(x^{\prime}, 0, y_{n}, \eta^{\prime}\right)=S_{2}\left(x^{\prime}, 0, y_{n}, \eta^{\prime}\right) \neq 0 \tag{1.18}
\end{equation*}
$$

on the other hand, the nondegeneracy of the blow-down implies that

$$
\begin{equation*}
\frac{\partial^{3} S\left(x^{\prime}, 0, y_{n}, \eta^{\prime}\right)}{\partial y_{n} \partial x_{n}^{2}}=\frac{\partial S_{2}}{\partial y_{n}}\left(x^{\prime}, 0, y_{n}, \eta^{\prime}\right) \neq 0 \tag{1.19}
\end{equation*}
$$

Conversely, one can easily show that any generating function of the form $S_{0}\left(x^{\prime}, \eta^{\prime}\right)+\frac{x_{n}^{2}}{2} S_{2}\left(x, y_{n}, \eta^{\prime}\right)$, with $S_{0}$ satisfying (1.15) and $S_{2}$ satisfying (1.18) and (1.19) gives rise to a FFCR. We have now proven

Theorem 1.20. $-A$ canonical relation $C \subset\left(T^{*} X \backslash 0\right) \times\left(T^{*} Y \backslash 0\right)$ is a fibered folding canonical relation near a point $\left(x_{0}, \xi_{0}, y_{0}, \eta_{0}\right)$, critical for $\pi$ (or $\rho$ ), iff there exist canonical transformations $\chi_{1}: T^{*} \mathbf{R}^{n} \backslash 0 \rightarrow T^{*} Y \backslash 0$, $\chi_{2}: T^{*} X \backslash 0 \rightarrow T^{*} \mathbf{R}^{n} \backslash 0$, with $\chi_{1}((0, \ldots, 0,1),(1,0, \ldots, 0))=\left(y_{0}, \eta_{0}\right)$, $\chi_{2}\left(x_{0}, \xi_{0}\right)=((0, \ldots, 0),(1,0, \ldots, 0,1))$, such that $\operatorname{Gr}\left(\chi_{2}\right) \circ C \circ \operatorname{Gr}\left(\chi_{1}\right)$ is parametrized by a phase function of the form

$$
\begin{equation*}
\phi\left(x, y, \eta^{\prime}\right)=S_{0}\left(x^{\prime}, \eta^{\prime}\right)-y^{\prime} \cdot \eta^{\prime}+\frac{x_{n}^{2}}{2} S_{2}\left(x, y_{n}, \eta^{\prime}\right) \tag{1.21}
\end{equation*}
$$

with $S_{0}$ and $S_{2}$ satisfying (1.15), (1.18) and (1.19).

## 2. $I^{p, \ell}(\Delta, \Lambda)$ and iterated regularity.

We now review the spaces of distributions associated with two cleanly intersecting Lagrangians [26], [16]; their characterization by means of iterated regularity [13]; and the $L^{2}$ estimates for operators whose Schwartz kernels are of this type [13]. Since only codimension 1 intersection is relevant to this paper, we will restrict our attention to that case. In the model case $\tilde{\Delta}=$ $\Delta_{T * \mathbf{R}^{n}}, \quad \tilde{\Lambda}=\left\{\left(x^{\prime}, x_{n}, \xi^{\prime}, 0 ; x^{\prime}, y_{n}, \xi^{\prime}, 0\right): x \in \mathbf{R}^{n}, \xi^{\prime} \in \mathbf{R}^{n-1} \backslash 0, y_{n} \in \mathbf{R}\right\}=$ the flowout of $\left\{\xi_{n}=0\right\}, I^{p, \ell}\left(\tilde{\Delta}^{\prime}, \tilde{\Lambda}^{\prime}\right)$ is defined to be the space of all sums of $C_{8}$ functions and distributions on $\mathbf{R}^{n} \times \mathbf{R}^{n}$ of the form

$$
\begin{equation*}
u(x, y)=\int e^{i\left(\left(x^{\prime}-y^{\prime}\right) \cdot \xi^{\prime}+\left(x_{n}-y_{n}-s\right) \cdot \xi_{n}+s . \sigma\right)} a(x, y, s ; \xi ; \sigma) d \sigma d s d \xi \tag{2.1}
\end{equation*}
$$

where $a$ is a product type symbol of order $p^{\prime}=p-\frac{n}{2}+\frac{1}{2}$, $\ell^{\prime}=\ell-\frac{1}{2}$, satisfying

$$
\begin{equation*}
\left|\partial_{\xi}^{\alpha} \partial_{\sigma}^{\beta} \partial_{x, y, s}^{\gamma} a\right| \leqslant C_{\alpha \beta \gamma K}(1+|\xi|)^{p^{\prime}-|\alpha|}(1+|\sigma|)^{\ell^{\prime \prime}-|\beta|} \tag{2.2}
\end{equation*}
$$

on each compact $K \subset \mathbf{R}_{x}^{n} \times \mathbf{R}_{y}^{n} \times \mathbf{R}_{s}$. In general, for a canonical relation $\Lambda \subset\left(T^{*} Y \backslash 0\right) \times\left(T^{*} Y \backslash 0\right)$ that intersects $\Delta_{T^{*} Y}$ cleanly in codimension 1 , one can find microlocally a canonical transformation $\chi$ : $\left(T^{*} Y \backslash 0\right) \times\left(T^{*} Y \backslash 0\right) \rightarrow\left(T^{*} \mathbf{R}^{n} \backslash 0\right) \times\left(\mathrm{T}^{*} \mathbf{R}^{n} \backslash 0\right)$ taking the pair $(\Delta, \Lambda)$ to $(\tilde{\Delta}, \tilde{\Lambda}) ; I^{p, \ell}\left(\Delta^{\prime}, \Lambda^{\prime}\right)$ is defined as the space of all microlocally finite sums of distributions $F_{i} u_{i}$, with $u_{i}$ of the form (2.1) and $F_{i} \in I^{0}\left(\operatorname{Gr}(\chi) ; \mathbf{R}^{n} \times \mathbf{R}^{n}, Y \times Y\right)$ for such a $\chi . I^{p, \ell}(\Delta, \Lambda)$ is then the class of operators with Schwartz kernel in $I^{p, \ell}\left(\Delta^{\prime}, \Lambda^{\prime}\right)$; microlocally if $T \in I^{p, \ell}(\Delta, \Lambda), T \in I^{p+\ell}(\Delta \backslash \Lambda ; Y)$ and $T \in I^{p}(\Lambda \backslash \Delta ; Y)$. Furthermore, the principal symbol of $T$ on $\Delta \backslash \Lambda$ lies in the space $R^{\bar{\ell}-\frac{1}{2}}$ defined in [16] and has a conormal singularity of order $\ell-\frac{1}{2}$ at $\Lambda$. The leading term of this singularity belongs to the space $S^{p, \ell}(Y \times Y ; \Delta, \Delta \cap \Lambda)$ of [16] and is denoted by $\sigma_{0}(T)$, the principal symbol of $T$ as an element of $I^{p, \ell}(\Delta, \Lambda)$.

The oscillatory representation (2.1) can be difficult to verify directly. Instead, we make use of the following characterization of $I^{p, t}\left(\Delta^{\prime}, \Lambda^{\prime}\right)$ from [13], which is a variant of the iterated regularity characterizations given by Melrose [22], [24] for various classes of distributions.

Proposition 2.3. - Let $\Lambda \subset\left(Y^{*} Y \backslash 0\right) \times\left(T^{*} Y \backslash 0\right)$ be a canonical relation cleanly intersecting the diagonal $\Delta$ in codimension 1. Then $u \in I^{p, \ell}\left(\Delta^{\prime}, \Lambda^{\prime}\right)$ for some $p, \ell \in \mathbf{R}$ iff for some $s_{0} \in \mathbf{R}$ and all $k \geqslant 0$, and all first order pseudodifferential operators $P_{1}\left(z, D_{z}, y, D_{y}\right)$, $P_{2}\left(z, D_{z}, y, D_{y}\right), \ldots$, whose principal symbols vanish on $\Delta^{\prime} \cup \Lambda^{\prime}$,

$$
\begin{equation*}
P_{1} \ldots P_{k} u \in H_{\mathrm{loc}}^{s_{0}}(Y \times Y) \tag{2.4}
\end{equation*}
$$

In the model case $(\tilde{\Delta}, \tilde{\Delta})$, the principal symbol of a first order $P\left(z, D_{z}, y, D_{y}\right)$, characteristic for $\tilde{\Delta}^{\prime} \cup \tilde{\Lambda}^{\prime}$, can be written (via the preparation theorem)

$$
\begin{equation*}
p(z, \zeta y, \eta)=\sum_{j=1}^{n} p_{j}\left(\zeta_{j}+\eta_{j}\right)+\sum_{j=1}^{n-1} q_{j}\left(z_{j}-y_{j}\right)+q_{n}\left(\zeta_{n}-\eta_{n}\right)\left(z_{n}-y_{n}\right) \tag{2.5}
\end{equation*}
$$

where the $p_{j}, q_{j}$ and $q_{n}$ are homogeneous of degrees 0,1 and 0 , respectively.

Finally, the following estimates are proven in [13], using the functional calculus of Antoniano and Uhlmann [1] and Jiang and Melrose (unpublished).

Theorem 2.6. - Let $\Sigma \subset T^{*} Y \backslash 0$ be a smooth, conic, codimension 1 submanifold and $\Lambda \subset\left(T^{*} Y \backslash 0\right) \times\left(T^{*} Y \backslash 0\right)$ its flowout. Then, if $T \in I^{p, \ell}(\Delta, \Lambda), T: H_{\text {comp }}^{s}(Y) \rightarrow H_{\text {loc }}^{s+s_{0}}(Y), \forall s \in \mathbf{R}$, if

$$
\begin{equation*}
\max \left(p+\frac{1}{2}, p+\ell\right) \leqslant s_{0} \tag{2.7}
\end{equation*}
$$

## 3. Composition and loss of $\frac{1}{4}$-derivative.

Let $C \subset\left(T^{*} X \backslash 0\right) \times\left(T^{*} Y \backslash 0\right)$ be a FFCR and $A \in I^{m}(C ; X, Y)$, $B \in I^{m^{\prime}}\left(C^{t} ; Y, X\right)$ properly supported FIOs.

Let $\Lambda=\Lambda_{\pi(L)}$ be the flowout of $\pi(L)$ in $\left(T^{*} Y \backslash 0\right) \times\left(T^{*} Y \backslash 0\right)$. By a microlocal partition of unity, we may write $A$ and $B$ as locally finite sums of operators $A=\sum A_{i}, B=\sum B_{j}$, such that on each $W F\left(A_{i}\right)^{\prime}$ or $W F\left(B_{j}\right)^{\prime}$, either $C$ is a canonical graph or Theorem 1.20 applies.

Furthermore, if $W F\left(B_{j}\right)^{\prime} \circ W F\left(A_{i}\right)^{\prime} \subset \Lambda$ (i.e., there is no contribution from the diagonal), then the clean intersection calculus of [9] and [32] applies, with excess $e=0$, to give $B_{j} A_{i} \in I^{m+m^{\prime}}(\Lambda ; Y, Y) \subset$
$I^{m+m^{\prime 0}}(\Delta, \Lambda ; Y, Y)$. We may thus restrict our attention to a composition $B A$, where $A \in I^{m}\left(C ; \mathbf{R}^{n}, \mathbf{R}^{n}\right), \quad B \in I^{m^{\prime}}\left(C^{t} ; \mathbf{R}^{n}, \mathbf{R}^{n}\right)$, with $C \subset\left(T^{*} \mathbf{R}^{n} \backslash 0\right) \times\left(T^{*} \mathbf{R}^{n} \backslash 0\right)$ parametrized by a phase function $\phi\left(x, y, \theta^{\prime}\right)=S_{0}\left(x^{\prime}, \theta^{\prime}\right)-y^{\prime} \cdot \theta^{\prime}+\frac{x_{n}^{2}}{2} S_{2}\left(x, y_{n}, \theta^{\prime}\right), \quad S_{0} \quad$ and $\quad S_{2} \quad$ satisfying (1.15), (1.18) and (1.19) in a conic neighborhood of $x=0, y_{n}=1$, $\theta^{\prime}=(1,0, \ldots, 0)$. By Hörmander's theorem [18], $A$ has an oscillatory representation

$$
\begin{equation*}
A f(x)=\int e^{i\left(S_{0}\left(x^{\prime}, \theta^{\prime}\right)-y^{\prime} \cdot \theta^{\prime}+\frac{x_{n}^{2}}{2} S_{2}\left(x, y_{n}, \theta^{\prime}\right)\right)} a\left(x, y, \theta^{\prime}\right) f(y) d \theta^{\prime} d y \tag{3.1}
\end{equation*}
$$

modulo a smoothing operator, where $a \in S_{1,0}^{m, \frac{1}{2}}\left(\mathbf{R}^{n} \times \mathbf{R}^{n} \times\left(\mathbf{R}^{n-1} \backslash 0\right)\right)$ is supported on a suitably small conic neighborhood of $x=(0, \ldots, 0)$, $y=(0, \ldots, 0,1), \theta^{\prime}=(1,0, \ldots, 0) . S_{0}\left(x^{\prime}, \theta^{\prime}\right)$ is, by (1.15), the generating function of a canonical transformation $\chi^{0}: T^{*} \mathbf{R}^{n-1} \backslash 0 \rightarrow T^{*} \mathbf{R}^{n-1} \backslash 0$, which we denote by $\left(\chi_{x^{\prime}}^{0}\left(x^{\prime}, \xi^{\prime}\right), \chi_{\xi^{\prime}}^{0}\left(x^{\prime}, \xi^{\prime}\right)\right)$; we may assume that $\chi^{0}\left(0, e_{1}^{*}\right)=\left(0, e_{1}^{*}\right)$. Then $\chi=\chi^{0} \otimes \mathrm{Id}: T^{*} \mathbf{R}^{n} \backslash 0 \rightarrow T^{*} \mathbf{R}^{n} \backslash 0$ is a canonical transformation. Let $F$ be a zeroth order FIO associated with $\chi^{-1}$, elliptic on $\rho(C) . F$ has the representation

$$
\begin{gathered}
F f(w)=\int e^{i\left(-S_{0}\left(x^{\prime}, \omega^{\prime}\right)+w^{\prime} \cdot \omega^{\prime}+\left(w_{n}-x_{n}\right) \cdot \omega_{n}\right)} c(x, w, \omega) f(x) d w d x \\
c \in S_{1,0}^{0}\left(\mathbf{R}^{n} \times \mathbf{R}^{n} \times\left(\mathbf{R}^{n} \backslash 0\right)\right)
\end{gathered}
$$

We compute the composition $F A$, applying as usual stationary phase in the $x, \omega$ variables. The critical points are given by $\omega^{\prime}=\theta^{\prime}+$ $\frac{x_{n}^{2}}{2} g_{1}\left(w^{\prime}, x_{n}, y_{n}, \theta^{\prime}\right), g_{1}$ smooth $\mathbf{R}^{n-1}$-valued and homogeneous of degree $1, \omega_{n}=0, x_{n}=w_{n}$, and $x^{\prime}$ determined by $w^{\prime}=d_{\omega^{\prime}} S_{0}\left(x^{\prime}, \omega^{\prime}\right)$, so that $x^{\prime}=\chi_{x^{\prime}}^{0}\left(w^{\prime}, \theta^{\prime}\right)+\frac{x_{n}^{2}}{2} g_{0}\left(w^{\prime}, x_{n}, y_{n}, \theta^{\prime}\right), g_{0}$ smooth and homogeneous of degree 0 . We thus have an oscillatory expression for $F A$ with symbol of order $m-\frac{1}{2}$ and phase

$$
\begin{equation*}
\left(w^{\prime}-y^{\prime}\right) \cdot \theta^{\prime}+\frac{w_{n}^{2}}{2}\left(S_{2}\left(x^{\prime}, w_{n}, y_{n}, \theta^{\prime}\right)+g_{1} \cdot\left(-d_{\theta^{\prime}} S_{0}\left(x^{\prime}, \theta^{\prime}\right)+w^{\prime}\right)\right) \tag{3.3}
\end{equation*}
$$

Since both $d_{\theta^{\prime}} S_{0}$ and $w^{\prime}$ vanish at $w=0, \quad y=(0, \ldots, 0,1)$, $\theta^{\prime}=(1,0, \ldots, 0)$, conditions (1.18) and (1.19) are still satisfied (if the
conic support of $A$ has been chosen suitably small to start with). Relabeling $w$ by $x$, one obtains

$$
\begin{equation*}
F A f(x)=\int e^{i\left(\left(x^{\prime}-y^{\prime}\right) \cdot \theta^{\prime}+\frac{x_{n}^{2}}{2} \tilde{S}_{2}\left(x, y_{n}, \theta^{\prime}\right)\right.} \tilde{a}\left(x, y, \theta^{\prime}\right) f(y) d \theta^{\prime} d y \tag{3.4}
\end{equation*}
$$

with $\tilde{S}_{2}$ satisfying (1.18) and (1.19) and $\tilde{a} \in S_{1,0}^{m-\frac{1}{2}}$, a refinement on the operator level of (1.21).
$F^{*} F$ is a zeroth order pseudodifferential operator $P$, elliptic on $\rho(C)$; let $Q$ be a property supported parametrix, so that $Q P=I \bmod C^{\infty}$ on distributions with wave-front set in $\rho(C)$. Then $B Q \in I^{m^{\prime}}\left(C^{t} ; \mathbf{R}^{n}, \mathbf{R}^{n}\right)$ and by repeating the above argument we obtain for $B Q F^{*}$ an oscillatory representation adjoint to (3.4), with symbol $\tilde{b} \in S_{1,0}^{m^{\prime}, \frac{1}{2}}$. Hence, modulo a smooth kernel, (cf. [8] [18]) the Schwartz kernel of BA has the following representation as an oscillatory integral :

$$
\begin{equation*}
K_{B A}(z, y)=\int e^{i\left(\left(x^{\prime}-y^{\prime}\right) \cdot \theta^{\prime}-\left(x^{\prime}-z^{\prime}\right) \sigma^{\prime}+\frac{x_{n}^{2}}{2}\left(\tilde{S}_{2}\left(x, y_{n}, \theta^{\prime}\right)-\tilde{S}_{2}\left(x, z_{n}, \sigma^{\prime}\right)\right)\right)} c d \theta^{\prime} d \sigma^{\prime} d x \tag{3.5}
\end{equation*}
$$

where $c \in S_{1,0}^{m+m^{\prime-1}}$ is $\tilde{a} \cdot \tilde{b}$ cutoff to be supported in $\left\{\left|\theta^{\prime}\right| \simeq\left|\sigma^{\prime}\right|\right\}$.
Now, since the gradient of the phase $\Phi\left(z, y, x, \theta^{\prime}, \sigma^{\prime}\right)=$ $\left(x^{\prime}-y^{\prime}\right) \cdot \theta^{\prime}-\left(x^{\prime}-z^{\prime}\right) \cdot \sigma^{\prime}+\frac{x_{n}^{2}}{2}\left(\tilde{S}_{2}\left(x, y_{n}, \theta^{\prime}\right)-\tilde{S}_{2}\left(x, z_{n}, \sigma^{\prime}\right)\right)$ in all the variables is $\neq 0$, integration by parts a finite number of times shows that all expressions of the form (3.5), with amplitude in $S_{1,0}^{m+m^{\prime-1}}$, lie in a fixed Sobolev space $H_{\text {loc }}^{s_{0}}\left(\mathbf{R}^{n} \times \mathbf{R}^{n}\right)$; in fact, we may take $s_{0}$ to be any number $<-\left(3 n+m+m^{\prime}-4\right)$ (cf., [18], p. 90).

Proposition 3.6. - For $x_{n}$ sufficiently small, there are smooth functions $C\left(y, z, x, \theta^{\prime}, \sigma^{\prime}\right)$ and $D\left(y, z, x, \theta^{\prime}, \sigma^{\prime}\right)$, taking values in $\mathbf{R}^{n}$ and $\operatorname{Hom}\left(\mathbf{R}^{n^{*}}, \mathbf{R}^{n-1}\right)$ and homogeneous of degrees -1 and 0 , respectively, such that

$$
\begin{equation*}
x_{n}\left(z_{n}-y_{n}\right) e^{i \Phi}=C \cdot d_{x}\left(e^{i \Phi}\right) \tag{3.7}
\end{equation*}
$$

and

$$
\begin{equation*}
\left(\sigma^{\prime}-\theta^{\prime}\right) e^{i \Phi}=D\left(d_{x} e^{i \Phi}\right) \tag{3.8}
\end{equation*}
$$

Proof. - Vanishing as it does at $\left\{z_{n}=y_{n}, \sigma^{\prime}=\theta^{\prime}\right\}, \quad \tilde{S}_{2}\left(x, y_{n}, \theta^{\prime}\right)$ $-\tilde{S}_{2}\left(x, z_{n}, \sigma^{\prime}\right)$ may be written as $\left(z_{n}-y_{n}\right) A\left(z, y, x, \theta^{\prime}, \sigma^{\prime}\right)$
$+B\left(z, y, x, \theta^{\prime}, \sigma^{\prime}\right) \cdot\left(\sigma^{\prime}-\theta^{\prime}\right)$, where $A$ and $B$ are smooth, $\mathbf{R}$-and $\mathbf{R}^{n-1}$ valued and homogeneous of degrees 1 and 0 , respectively. By (1.19), $A \neq 0$ near $z=y, x_{n}=0, \theta^{\prime}=\sigma^{\prime}$. Then we have

$$
\begin{equation*}
d_{x_{n}} \Phi=x_{n}\left(\left(z_{n}-y_{n}\right)\left(A+\frac{x_{n}}{2} d_{x_{n}} A\right)+\left(\sigma^{\prime}-\theta^{\prime}\right) \cdot\left(B+\frac{x_{n}}{2} d_{x_{n}} B\right)\right) \tag{3.9}
\end{equation*}
$$

and

$$
\begin{equation*}
d_{x^{\prime}} \Phi=\theta^{\prime}-\sigma^{\prime}+\frac{x_{n}^{2}}{2}\left(\left(z_{n}-y_{n}\right) d_{x^{\prime}} A+\left(\sigma^{\prime}-\theta^{\prime}\right) \cdot d_{x^{\prime}} B\right) \tag{3.10}
\end{equation*}
$$

Solving (3.10), we have

$$
\begin{equation*}
\left(I-\frac{x_{n}^{2}}{2} d_{x^{\prime}} B\right)\left(\sigma^{\prime}-\theta^{\prime}\right)=-d_{x^{\prime}} \Phi+\frac{x_{n}^{2}}{2}\left(z_{n}-y_{n}\right) d_{x^{\prime}} A \tag{3.11}
\end{equation*}
$$

and combining this with (3.9) we have, for $x_{n}$ small,
(3.12) $\quad x_{n}\left(z_{n}-y_{n}\right)$

$$
=\frac{1}{\tilde{A}}\left(x_{n}\left(I-\frac{x_{n}^{2}}{2} d_{x^{\prime}} B\right)^{-1^{*}}\left(B+\frac{x_{n}^{2}}{2} d_{x_{n}} B\right) \cdot d_{x^{\prime}} \Phi+d_{x_{n}} \Phi\right)
$$

where

$$
\tilde{A}=A+\frac{x_{n}^{2}}{2} d_{x_{n}} A+\frac{x_{n}^{2}}{2}\left(I-\frac{x_{n}^{2}}{2} d_{x^{\prime}} B\right)^{-1 *}\left(B+\frac{x_{n}}{2} d_{x_{n}} B\right) \cdot d_{x^{\prime}} A \neq 0
$$

implying (3.7). From this and the step following (3.11) we obtain (3.8).
We are now in a position to verify that $K_{B A} \in I^{p, \ell}\left(\Delta^{\prime}, \Lambda^{\prime}\right)$, for some $p, \ell \in \mathbf{R}$, using iterated regularity. Given a first order $P\left(z, D_{z}, y, D_{y}\right)$, characteristic for $\Delta^{\prime} \cup \Lambda^{\prime}$, we recall from (2.5) that its principal symbol may be written

$$
p(z, \zeta, y, \eta)=\sum_{1}^{n} p_{j}\left(\zeta_{j}+\eta_{j}\right)+\sum_{1}^{n-1} q_{j}\left(z_{j}-y_{j}\right)+q_{n}\left(z_{n}-y_{n}\right)\left(\zeta_{n}-\eta_{n}\right)
$$

By (3.5), we have (cf. [8])
(3.13) $P K_{B A}(z, y)=\int e^{i \varphi\left(z, y, x, \theta^{\prime}, \sigma^{\prime}\right)}\left(p\left(z, d_{z} \Phi, y, d_{y} \Phi\right) c+d\right) d \theta^{\prime} d \sigma^{\prime} d x$,
with $d \in S_{1,0}^{m+m^{\prime}-1}$. Since $d_{z^{\prime}} \Phi+d_{y^{\prime}} \Phi=\sigma^{\prime}-\theta^{\prime}$, if we let $p^{\prime}=\left(p_{1}, \ldots, p_{n-1}\right)$, the $p^{\prime} \cdot\left(\zeta^{\prime}+\eta^{\prime}\right)$ term of $P K_{B A}$ is

$$
\begin{aligned}
\int e^{i \Phi} p^{\prime} \cdot\left(\sigma^{\prime}-\theta^{\prime}\right) c d \theta^{\prime} d \sigma^{\prime} d x & =\int D\left(d_{x} e^{i \Phi}\right) \cdot p^{\prime} c d \theta^{\prime} d \sigma^{\prime} d x \\
& =\int e^{i \Phi} d_{x}^{t} D^{*}\left(p^{\prime} c\right) d \theta^{\prime} d \sigma^{\prime} d x
\end{aligned}
$$

by (3.8); but because $D$ is homogeneous of degree 0 , $d_{x}^{t} D^{*}\left(p^{\prime} c\right) \in S_{1,0}^{m+m^{\prime}-1}$ and this is of the form (3.5). For the $p_{n}\left(\zeta_{n}+\eta_{n}\right)$ term, note that

$$
\left.d_{z_{n}} \Phi+d_{y_{n}} \Phi=\frac{x_{n}^{2}}{2}\left(\left(z_{n}-y_{n}\right) d_{z_{n}} A+d_{y_{n}} A\right)+\left(\sigma^{\prime}-\theta^{\prime}\right) \cdot\left(d_{z_{n}} B+d_{y_{n}} B\right)\right),
$$

leading to
$\int e^{i \Phi} d_{x}^{t} \cdot\left(C^{*}\left(\frac{x_{n} p_{n} c}{2}\left(d_{z_{n}} \mathrm{~A}+d_{y_{n}} A\right)\right)+D^{*}\left(p_{n} c\left(d_{y_{z_{n}}} B+d_{y_{n}} B\right)\right)\right) d \theta^{\prime} d \sigma^{\prime} d x$,
which is again of the form (3.5). Similarly, noting $\left.d_{\sigma^{\prime}} \Phi+d_{\theta^{\prime}} \Phi=z^{\prime}-y^{\prime}+\frac{x_{n}^{2}}{2}\left(\left(z_{n}-y_{n}\right)\left(d_{\sigma^{\prime}} A+d_{\theta^{\prime}} A\right)+\left(\sigma^{\prime}-\theta^{\prime}\right) \cdot d_{\sigma^{\prime}} B+d_{\theta^{\prime}} B\right)\right)$, we find that

$$
\begin{align*}
\left(z^{\prime}-y^{\prime}\right) e^{i \Phi}=i^{-1}\left(d_{\sigma^{\prime}}+d_{\theta^{\prime}}\right) e^{i \Phi} & -\frac{x_{n}}{2}\left(d_{\sigma^{\prime}} A+d_{\theta^{\prime}} A\right) C \cdot d_{x} e^{i \Phi}  \tag{3.14}\\
& -\frac{x_{n}^{2}}{2} D^{*}\left(d_{\sigma^{\prime}} B+d_{\theta^{\prime}} B\right) \cdot d_{x} e^{i \Phi}
\end{align*}
$$

and thus the $\sum_{1}^{n-1} q_{j}\left(z_{j}-y_{j}\right)$ term of $P K_{B A}$ is of the form (3.5). Finally,

$$
\begin{aligned}
d_{z_{n}} \Phi & -d_{y_{n}} \Phi \\
& =x_{n}\left(x_{n}\left(x_{n} A+\frac{x_{n}}{2}\left(z_{n}-y_{n}\right) d_{z_{n}} A-d_{y_{n}} A\right)+\frac{x_{n}}{2}\left(\sigma^{\prime}-\theta^{\prime}\right) \cdot\left(d_{z_{n}} B-d_{y_{n}} B\right)\right),
\end{aligned}
$$

so that the $q_{n}\left(z_{n}-y_{n}\right)\left(\zeta_{n}-\eta_{n}\right)$ term of $P K_{B A}$ is

$$
\int e^{i \Phi} d_{x}^{t} \cdot C^{*}\left(x_{n} A+\ldots\right) d \theta^{\prime} d_{\sigma^{\prime}} d x
$$

again an oscillatory integral of the form (3.5) with symbol in $S_{1,0}^{m+m^{\prime}-1}$. By induction, for any first order operators $P_{1}, \ldots, P_{k}$, characteristic for $\Delta^{\prime} \cup \Lambda^{\prime}, P_{1}, \ldots, P_{k} K_{B A}$ is of this form, and hence in $H_{\text {loc }}^{s_{0}}\left(\mathbf{R}^{n} \times \mathbf{R}^{n}\right)$ by the comment above.

Prop. 2.3 yields $K_{B A} \in I^{p, \ell}\left(\Delta^{\prime}, \Lambda^{\prime}\right)$ and hence $B A \in I^{p, \ell}(\Delta, \Lambda)$, for some $p, \ell \in \mathbf{R}$.

To determine the orders $p$ and $\ell$, note that away from $L$ the composition is covered by Hörmander's calculus and hence $B A \in I^{m+m^{\prime}}(\Delta \backslash \Lambda ; Y, Y)$ microlocally so that $p+\ell=m+m^{\prime}$. Furthermore, the calculation of the principal symbol of $B A$ in [18] is still valid away from $\pi(L)$. If $a$ is the principal symbol of $A$, considered as a $\frac{1}{2}$-density on $C$, we may express $a$ as $\alpha \cdot\left|\pi^{*} \omega_{Y}^{n}\right|^{1 / 2}$. Since $\pi^{*} \omega_{Y}=\omega_{C}$ is folded sympletic, $\pi^{*} \omega_{Y}^{n}$ vanishes to first order at $L$ and thus $\alpha$ has a conormal singularity of order $-\frac{1}{2}$ at $L$.

Similarly, the principal symbol of $B$ is $b=\beta \cdot \mid \pi^{*} \omega_{Y^{\prime} \mid{ }^{1 / 2}}$ with $\beta$ having a conormal singularity of order $-\frac{1}{2}$ at $L^{t}$ (here $Y^{\prime}$ denotes the second copy of $Y$ ). Thus $\left.\beta \cdot \alpha\right|_{T_{*} Y^{\prime} \times \Delta_{T_{*}} \times{ }^{\times}{ }^{*} Y_{Y}}$ has a conormal singularity of order -1 above $\pi(L)$; when pushed down by the Whitney fold $\pi$, this gives rise to a conormal singularity of order $-\frac{1}{2}$ at $L$, in the principal symbol $b \times a$ of $B A$ (cf. [12]). Hence, $\ell-\frac{1}{2}=-\frac{1}{2}$, and $p=m+m^{\prime}, \ell=0$, finishing the proof of Theorem 0.1. In addition, we see that the principal symbol $\sigma_{0}(B A)$ is the image of $b \times a$ in $S^{m+m^{\prime}, 0}(Y \times Y ; \Delta, \pi(L))$.

To prove Corollary 0.2 , suppose $A \in I^{m}(C ; X, Y)$ is properly supported, with $C$ a FFCR. Then $A^{*} A \in I^{2 m, 0}\left(\Delta, \Lambda_{\pi(L)} ; Y, Y\right)$ and is properly supported and so maps $H_{\text {comp }}^{s}(Y) \rightarrow H_{\text {loc }}^{s-2 m-1 / 2}(Y)$ by Theorem 2.6. This yields Corollary 0.2 for $s=m+\frac{1}{4}$. For general $s \in \mathbf{R}$, we simply apply this result to PAQ, where $P$ and $Q$ are elliptic pseudodifferential operators on $X$ and $Y$ of orders $-s+m+\frac{1}{4}$ and $s-m-\frac{1}{4}$,
respectively. As shown by an example in [13], one does not lose less than $\frac{1}{4}$ derivative in general.

It is also possible to give sharp estimates for $A$ in terms of nonisotropic Sobolev spaces. Let $\Psi^{m}(Z)$ denote the pseudodifferential operators of order $m$ and type 1,0 on a manifold $Z$. Then, for $s \in \mathbf{R}$,

$$
\begin{align*}
H_{\mathrm{loc}}^{s, k}(X)=\left\{v \in \mathscr{D}^{\prime}(X):\right. & Q_{1} \ldots Q_{k} v \in H_{\mathrm{loc}}^{s}(X)  \tag{3.15}\\
& \text { for all } \left.Q_{j} \in \Psi^{1}(X) \text { with }\left.\sigma_{\mathrm{prin}}\left(Q_{j}\right)\right|_{\mathrm{p}(L)}=0, \forall j\right\}
\end{align*}
$$

is the nonisotropic Sobolev space of [3]; defined initially for $k \in \mathbf{Z}_{+}$, one uses interpolation and duality to extend the definition to $k \in \mathbf{R}$. Since $\rho(L)$ is symplectic, we have $H_{\mathrm{loc}}^{s, k}(X) \hookrightarrow H_{\mathrm{loc}}^{s+k / 2}(X)$; microlocally away from $\rho(L)$, of course, $H_{\mathrm{loc}}^{s, k}(X) \hookrightarrow H_{\mathrm{loc}}^{s+k}(X)$. For $s \in \mathbf{R}$, set

$$
\begin{align*}
& H_{\mathrm{loc}}^{s, k}(Y)=\left\{u \in \mathscr{D}^{\prime}(Y): P_{1} \ldots P_{k} u \in H_{\mathrm{loc}}^{s}(Y)\right.  \tag{3.16}\\
& \\
& \text { for all } \left.P_{j} \in \Psi^{1}(Y) \text { with }\left.\sigma_{\text {prin }}\left(P_{j}\right)\right|_{\pi(L)}=0, \forall j\right\},
\end{align*}
$$

again extended to $k \in \mathbf{R}$ by interpolation and duality. (For $\pi(L)$ the characteristic variety of the wave operator, this space has been widely used in the study of nonlinear problems.) One can then show that if $A \in I^{m}(C ; X, Y)$ is properly supported, with $C$ a FFCR,

$$
\begin{equation*}
A: H_{\mathrm{loc}}^{s, k}(Y) \rightarrow H_{\mathrm{loc}}^{s-k-m-1 / 2,2 k+1 / 2}(X), \tag{3.17}
\end{equation*}
$$

giving a sharper form of (0.2). The main point in the proof is to show that if $Q_{1}, Q_{2}, \in \Psi^{1}(X)$ are characteristic for $\rho(L)$, then there are operators $P_{1}, P_{2} \in \Psi^{1}(Y)$ characteristic for $\pi(L)$ and $A_{1}, A_{2}$, $A_{3} \in I^{m+1}(C ; X, Y)$ such that $Q_{1} Q_{2} A=A_{1} P_{1}+A_{2} P_{2}+A_{3}$. This is done by splitting $\rho^{*}\left(\sigma_{\text {prin }}\left(Q_{1}\right) \sigma_{\text {prin }}\left(Q_{2}\right)\right)$ into its even and odd components with respect to the fold involution of $C$. The details are left to the reader.

## 4. $\boldsymbol{L}^{\boldsymbol{p}}$ estimates for restricted $\boldsymbol{X}$-ray transforms.

Let $(M, g)$ be an $n$-dimensional riemannian manifold. The hamiltonian function $H(x, \xi)=g(x, \xi)^{1 / 2}$ generates the geodesic flow on $T^{*} M \backslash 0$, which preserves $S^{*} M=\{(x, \xi): H(x, \xi)=1\}$. Suppose $M$ is such that
$S^{*} M$ modded out by this flow is a smooth, ( $2 n-2$ )-dimensional manifold, $\mathscr{M}$. This holds, for example, if the action of $\mathbf{R}$ on $S^{*} M$ given by the geodesic flow is free and proper, as is the case if $M$ is geodesically convex (e.g., $\mathbf{R}^{n}$ with the standard metric). $\mathscr{M}$ is also smooth if $M$ is a compact, rank one symmetric space [2]. One identifies $\mathscr{M}$ with the space of oriented geodesics on $M$ and then defines the $X$-ray transform (cf. Helgason [27])

$$
\begin{equation*}
\mathscr{R} f(\gamma)=\int_{\mathbf{R}} f(\gamma(s)) d s, \quad f \in C_{0}^{\infty}(M), \gamma \in \mathscr{M} \tag{4.1}
\end{equation*}
$$

where $\gamma(s)$ is any unit-velocity parametrization of $\gamma . \mathscr{R}$ is a generalized Radon transform in the sense of Guillemin, satisfying the Bolker condition, and hence the clean intersection calculus applies, yielding that $\mathscr{R}^{*} \mathscr{R}$ is a pseudodifferential operator of order -1 on $M$ [14]. Thus, $\mathscr{R}: H_{\text {comp }}^{s}(M) \rightarrow H_{l o c}^{s+1 / 2}(\mathscr{M})$, generalizing (locally) the result of Smith and Solmon [28] for the $X$-ray transform in $\mathbf{R}^{n}$.

One now considers the restriction of $\mathscr{R} f$ to $n$-dimensional submanifolds (geodesic complexes) $\mathscr{C} \subset \mathscr{M}$, and the question of reconstructing $f$ from $\mathscr{R}_{\mathscr{C}} f=\left.\mathscr{R} f\right|_{\mathscr{8}}$. (The following is a summary of the discussion in [12], to which the reader is referred for more details.) To even define $\mathscr{R}_{6} f$ for $f \in \mathscr{E}^{\prime}(M)$, we have to impose a restriction on the wave-front set of $f$. Let

$$
\begin{equation*}
Z_{\mathscr{C}}=\{(\gamma, x) \in \mathscr{C} M: x \in \gamma\} \tag{4.2}
\end{equation*}
$$

be the point-geodesic relation of $\mathscr{C}$; the Schwartz kernel of $\mathscr{R}_{8}$ is a smooth multiple of the delta function on $Z_{\mathscr{8}}$. Let Crit ( $\left.\mathscr{C}\right)$ be the critical values of the projection from $Z_{母}$ to $M$; by Sard's theorem, this is nowhere dense and of measure 0 . There is a closed conic set $K_{0} \subset T^{*} M \backslash 0$, whose complement sits over Crit ( $\mathscr{C}$ ), such that for

$$
f \in \mathscr{E}_{K_{0}}^{\prime}(M)=\left\{f \in \mathscr{E}^{\prime}(M): W F(f) \subset K_{0}\right\}, \mathscr{R}_{\mathscr{E}} f \in \mathscr{D}(\mathscr{C})
$$

is well-defined. Shrinking $K_{0}$ to a somewhat smaller $K$ in order to avoid the nonfold critical points of $\pi: C=N^{*} Z_{\&}^{\prime} \rightarrow T^{*} M \backslash 0$, in [12] it was shown that if $\mathscr{C}$ satisfies a generalization of Gelfand's admissibility criterion [11], then, over $K, C$ is a FFCR and we have $\mathscr{R}_{\mathscr{6}} \in I^{-1 / 2}(C ; \mathscr{C}, M)$. Using an explicit description of the integral kernel of $\mathscr{R}_{\mathscr{6}}^{*} \mathscr{R}_{\mathscr{G}}$, it was also shown that $\mathscr{R}_{\mathscr{6}}^{*} \mathscr{R}_{\mathscr{G}} \in I^{-1,0}\left(\Delta_{T_{*} M}, \Lambda_{\pi(L)}\right)$, where $\pi(L)$
is the boundary of the support of the Crofton symbol, allowing the construction of a relative left-parametrix for $\mathscr{R}_{\&}$. From Theorem 2.6 it then followed that

$$
\begin{equation*}
\left\|\mathscr{R}_{\mathscr{G}} f\right\|_{H^{s+1 / 4}(\mathscr{G})} \leqslant C_{s}\|f\|_{H^{s}(M)}, f \in \mathscr{E}_{K}^{\prime} s \geqslant-\frac{1}{4}, \tag{4.3}
\end{equation*}
$$

$C_{s}$ depending on $s$ and the support of $f$. It now follows directly from (0.2) that (4.3) holds for all $s \in \mathbf{R}$; furthermore, by (3.17), $\mathscr{R}_{\mathscr{6}}: H_{\mathrm{loc}}^{s, k}(M) \rightarrow H_{\text {loc }}^{s-k+1 / 4,2 k}(\mathscr{C})$. Moreover, ( 0.2 ) can be applied to an analytic continuation of $\mathscr{R}_{\mathscr{C}}$ to obtain Theorem 0.4.

First, we derive necessary conditions for local boundedness

$$
\begin{equation*}
\mathscr{R}_{\mathscr{C}}: L_{\text {comp }}^{p}(M) \rightarrow L_{\mathrm{loc}}^{q}(\mathscr{C}) \tag{4.4}
\end{equation*}
$$

by considering, in $\mathbf{R}^{n}$, the following two families of functions. If $x \in \mathbf{R}^{n} \backslash$ Crit ( $\mathscr{C}$ ), i.e., the projection from $Z_{\mathscr{G}}$ to $\mathbf{R}^{n}$ is a submersion at $x_{0}$, then if we set $f_{\varepsilon}=\chi_{B\left(x_{0} ; \varepsilon\right)}$, we have $\left\|f_{\varepsilon}\right\|_{L^{p}} \sim \varepsilon^{n / p}$ while $\mathscr{R}_{\varepsilon} f_{\varepsilon} \geqslant c \varepsilon$ on a rectangle in $\mathscr{C}$ of dimensions $\sim 1 \times \varepsilon \times \varepsilon^{n-1}$, so that $\left\|\mathscr{R}_{\varepsilon_{\varepsilon}} f_{\varepsilon}\right\|_{L^{q}} \geqslant c \varepsilon^{1+\frac{n-1}{q}} ;(4.4)$ then implies that $\frac{1}{q} \geqslant(n / n-1) \frac{1}{p}-\frac{1}{n-1}$. If $0=x_{0} \in \gamma_{0}=x_{1}-$ axis and $T_{\gamma_{0}} \Sigma=x_{1}-x_{2}$ plane, where

$$
\sum_{x_{0}}=\bigcup_{\left\{\gamma \in \mathscr{Q}: x_{0} \in \gamma\right\}}
$$

is a two-dimensional cone with vertex at $x_{0}$ and $T_{\gamma_{0}} \sum_{x_{0}}$ is its tangent plane along $\gamma_{0}$, we may set $f_{\varepsilon}=\chi_{[-1,1] \times[-\varepsilon, \varepsilon] \times\left[-\varepsilon^{2}, \varepsilon^{2}\right] \times \ldots \times\left[-\varepsilon^{2}, \varepsilon^{2}\right]}$, obtaining $\left\|f_{\varepsilon}\right\|_{L^{p}} \sim \varepsilon \frac{2 n-3}{p}$ while $\left\|\mathscr{R}_{\varepsilon} f\right\|_{L^{q}} \geqslant c \varepsilon \frac{2 n-2}{q}$, so that (4.4) implies that $\frac{1}{q} \geqslant(2 n-3) /(2 n-2) \cdot \frac{1}{q}$. Thus, a necessary condition for (4.4) to hold is that $\left(\frac{1}{p}, \frac{1}{q}\right)$ lie in the convex hull of $(0,0),(1,1)$ and $\left(\frac{2}{3},(2 n-3) /(3 n-3)\right)$. Our positive results, $(0.4 a)$ and $(0.4 b)$, are only for $\left(\frac{1}{p}, \frac{1}{q}\right)$ lying in a proper subset of this region and so are probably not sharp.

The proof of Theorem 0.4 is straightforward, given Theorem 0.2. Let $\rho_{1}(\gamma, x), \ldots, \rho_{n-1}(\gamma x) \in \mathscr{C}^{\infty}(\mathscr{C} \times M)$ be defining functions for $Z_{8}$. Consider the entire, distribution-valued family

$$
\begin{equation*}
K^{\alpha}(\gamma, x)=\Gamma\left(\frac{\alpha}{2}\right)^{-1}|\vec{\rho}(\gamma, x)|^{\alpha-(n-1)} \psi(\gamma, x), \alpha \in \mathbf{C} \tag{4.5}
\end{equation*}
$$

where $\vec{\rho}=\left(\rho_{1}, \ldots, \rho_{n-1}\right)$ and $\psi \in C_{0}^{\infty}(\mathscr{C} \times M)$ is $\equiv 1$ on $Z_{8}$ over the support of $f$ and supported close to $Z_{\mathscr{C}}$. If we denote the operator with Schwartz kernel $K^{\alpha}$ by $\mathscr{R}^{\alpha}$, then $\mathscr{R}^{\alpha} \in I^{-1 / 2-\operatorname{Re}(\alpha)}(C ; \mathscr{C}, M)$. Furthermore, if $P(x, D)$ is a zeroth order pseudodifferential operator on $M$, elliptic on a subcone $K_{1} \subset K$ and smoothing outside of $K$, then $\mathscr{R}^{0}=\mathscr{R}_{8} P$ acting on $\mathscr{E}_{K_{1}}^{\prime}$. By (0.2), we have $\mathscr{R}^{\alpha} P: L_{\text {comp }}^{2}(M) \rightarrow$ $L_{\text {loc }}^{2}(\mathscr{C})$ for $\operatorname{Re}(\alpha)=-\frac{1}{4}$. On the other hand, for $\operatorname{Re}(\alpha)=n-1$, we clearly have $\mathscr{R}^{\alpha} P: H^{1} \rightarrow L_{\text {loc }}^{\infty}$, where $H^{1}$ is the Hardy space on $M$ ([29]). By the Fefferman-Stein interpolation theorem [10],

$$
\mathscr{R}^{0}: L_{\mathrm{comp}}^{p_{0}} \rightarrow L_{\mathrm{comp}}^{q_{0}}\left(\frac{1}{p_{0}}, \frac{1}{q_{0}}\right)=\left(\frac{2 n-1}{4 n-3}, \frac{2 n-2}{4 n-3}\right) .
$$

(A word is needed about the dependence of the $L^{2}$ bounds on $\operatorname{Im}(\alpha)$ for $\operatorname{Re}(\alpha)=-\frac{1}{4}$. To obtain estimates on any finite number of derivatives of the product-type symbol of $\mathscr{R}^{\alpha^{*}} \mathscr{R}^{\alpha} \in I^{-1-2 \operatorname{Re}(\alpha), 0}(\Delta, \Lambda)$, only a finite number of applications of first order pseudodifferential operators (as in (2.3)) have to be made. However, the dependence of $L^{2}$ bounds for elements of $I^{-1 / 2,0}(\Delta, \Lambda)$ on only a finite number of derivatives of the product-type symbols is not clear in the proof presented in [13], § 3, since that proof uses the full functional calculus for $I(\Delta, \Lambda)$. An alternate proof may be given, though, in which this dependence is clear. There are fixed elliptic FIOs $F_{1}, F_{2}$ such that $T^{\alpha}=F_{2} \mathscr{R}^{\alpha *} \mathscr{R}^{\alpha} F_{1} \in I^{-1 / 2,0}(\tilde{\Delta}, \tilde{\Lambda})$ has the representation (cf. [13], § 1).

$$
T^{\alpha} f(z)=\int e^{i\left(\left(z^{\prime}-y^{\prime}\right) \cdot \zeta^{\prime}+\left(z_{n}-y_{n}\right) \zeta_{n}\right)} a_{\alpha}\left(z, y ; \zeta^{\prime} ; \zeta_{n}\right) f\left(y^{\prime}, y_{n}\right) d \zeta^{\prime} d \zeta_{n} d y^{\prime} d y_{n}
$$

where $a_{\alpha}$ is a symbol-valued symbol of order $M=0, M^{\prime}=0$. We may consider this as a pseudodifferential operator, of order 0 and type 1,0 , acting on $L^{2}\left(\mathbf{R}^{n-1} ;\left(L^{2}(\mathbf{R})\right)\right.$, whose symbol is the pseudodifferential operator on $\mathbf{R}$ with symbol $a_{\alpha}\left(z^{\prime}, \cdot, y^{\prime}, \cdot ; \zeta^{\prime} ; \cdot\right)$, which is of order 0 and type 1,0 . By the standard proofs of $L^{2}$ boundedness for operators
of type 1,0 , we only need the $S_{1,0}^{0}$ estimates for a finite number (say, $n$ ) of derivatives. Thus, the $L^{2}$ bounds for $\mathscr{R}^{\alpha}$ grow at most exponentially in $|\operatorname{Im}(\alpha)|$ for $\operatorname{Re}(\alpha)=-\frac{1}{4}$ ).

On compact sets away from Crit $(\mathscr{C}), \quad \sup _{x}\left\|K_{\mathscr{M}_{\mathscr{C}}}(\cdot, x)\right\|$ and $\sup _{\gamma}\left\|K_{\pi_{\varnothing}}(\gamma, \cdot)\right\|$ are bounded, where $\|d \mu\|$ is the total variation of a complex measure $d \mu$, and hence $\mathscr{R}_{\mathscr{C}}: L_{\text {comp }}^{p} \rightarrow L_{\text {ioc }}^{p} 1 \leqslant p \leqslant \infty$, acting on functions supported away from $\operatorname{Crit}(\mathscr{C})$, and hence $\mathscr{R}_{\mathscr{C}} P: L_{\text {comp }}^{p} \rightarrow$ $L_{\text {loc }}^{p} 1<p \leqslant \infty$. Interpolating between these estimates, we obtain Theorem 0.4. Of course, if we can take $K=T^{*} M \backslash 0$, then the microlocalization $P(x, D)$ is unnecessary and ( 0.4 ) holds for $p=1$, $p=\infty$ as well.

Just as with the $L^{2}$ estimates in [13], one expects the estimates for $\mathscr{R}_{8}$ for a general $\mathscr{C}$ to be better than those in (0.4). For instance, it was shown in [13] that for an open set of $\mathscr{C}$ 's in three variables, $N^{*} Z_{\mathscr{G}}^{\prime}$ is a folding canonical relation in the sense of Melrose and Taylor [25], so that there is a loss of only $\frac{1}{6}$, rather than $\frac{1}{4}$, derivatives on $L^{2}$. Incorporating the $L^{2}$ estimates of [25] into the above interpolation argument, one obtains

Theorem 4.6. - Let $\mathscr{C} \subset \mathscr{M}$ be a geodesic complex and let $P(x, D)$ be a zeroth order pseudodifferential operator on $M$ such that $C=N^{*} Z_{\mathscr{G}}^{\prime}$ is a folding canonical relation over the conic support of $P$. Then $\mathscr{R}_{\mathscr{G}} P: L_{\text {comp }}^{p}(M) \rightarrow L_{\text {loc }}^{q}(\mathcal{M})$ for $p, q$ satisfying either of the following conditions :
(a) $\frac{1}{q} \geqslant \frac{3 n-1}{3 n-3}\left(\frac{1}{p}-\frac{1}{2(3 n-2)}\right), \quad 1<p \leqslant \frac{2(3 n-2)}{3 n-1}$;
(b) $\frac{1}{q} \geqslant \frac{3 n-3}{3 n-1} \frac{1}{p}, \quad \frac{2(3 n-2)}{3 n-1} \leqslant p<\infty$.

As described in [13], examples of $\mathscr{C}$ 's to which Theorem 4.6 applies are given by equipping $\mathbf{R}^{3}$ with the Heisenberg group structure with Planck's constant $\varepsilon \neq 0$ suitably small and taking $\mathscr{C}_{\varepsilon}$ to be all light rays through the origin and their left translates. Because of the stability of Whitney folds, Theorem 4.6 also applies to small perturbations of these in the $C^{\infty}$ topology.

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